

Kepler Kolloquium

“Low-Energy Neutrino Detection through Neutrino-Nucleus Interaction Probes”

T.S. Kosmas

Division of Theoretical Physics, **University of Ioannina**, GR-45110, Greece

- **Low-energy Neutrino production Sources**
- **Neutrino Detection**
- **Neutrino-Nucleus Interactions**
 - **Coherent Neutrino-Nucleus Scattering**
 - **Inelastic Neutrino-Nucleus Scattering**

Collaborating Groups:

University of Ioannina, Greece : V. Tsikoudi, J. Sinatkas, V. Stavrou
K. Balasi, V. Tsakstara, G. Karathanou,
P. Giannaka, J. Kardaras
Hellenic Army Academy, Greece: P.C. Divari, Th. Liolios

Univ. of Tuebingen, Germany : Group of A. Faessler

T. Univ. of Darmstadt, Germany : Group of J. Wambach

Univ. of Jyvaskyla, Finland : Group of J. Suhonen

Univ. of Valencia, Spain : Group of J.W.F. Valle

RCNP, Univ. of Osaka, Japan : H. Ejiri (MOON-experiment)

T.Univ Dresden, Germany : K. Zuber (COBRA-experiment)

Important neutrino *production* sources

1) Laboratory Neutrino Sources

- i. Reactor neutrinos (Low-Energy neutrino beams)
 - Slow – pion and muon decay ($E < 52.8 \text{ MeV}$)
 - Beta decay ($E < 10 \text{ MeV}$)
- ii. Accelerator neutrinos (High-Energy neutrino beams)
 - Boosted beta-beam neutrinos ($E < 100 \text{ MeV}$ and more)
 - Long (short) baseline neutrinos

2) Astrophysical Neutrino Sources

- i. Solar Neutrinos ($E < 20 \text{ MeV}$)
- ii. Supernova Neutrinos ($E < 60 - 80 \text{ MeV}$)
- iii. Atmospheric Neutrinos ($E < \text{a few GeV}$)
- iv. Gamma-Ray-burst neutrinos, Cosmological Neutrinos, etc.

Reactor neutrino sources

Reactor neutrino beams used to study low-energy neutrino interactions and neutrino oscillations (KARMEN, LSND, BooNE, etc.):

slow-pion decay

$$\pi^+ \rightarrow \mu^+ + \nu_\mu$$

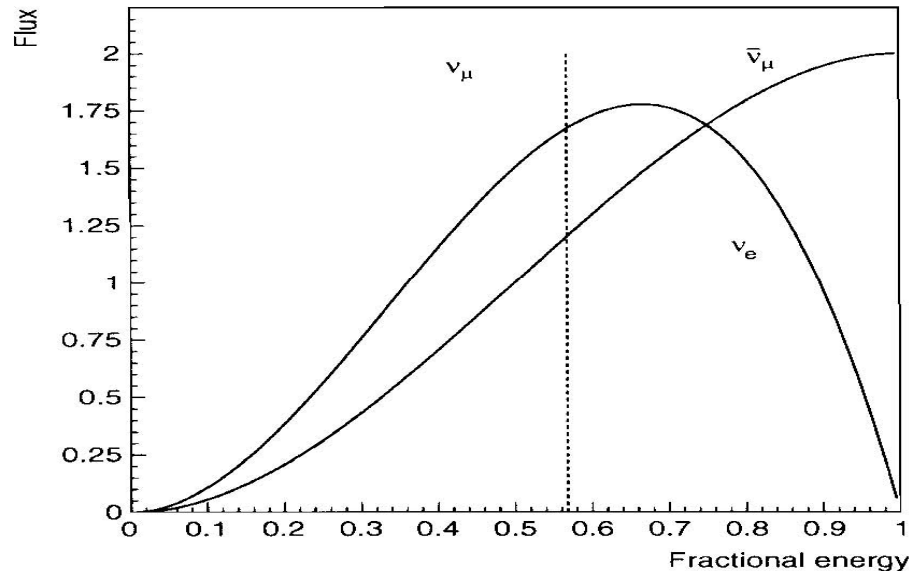
$$\pi^- \rightarrow \mu^- + \bar{\nu}_\mu$$

μ -decay (at rest)

$$\mu^+ \rightarrow e^+ + \nu_e + \bar{\nu}_\mu$$

$$\mu^- \rightarrow e^- + \bar{\nu}_e + \nu_\mu$$

W.C. Louis / Progress in Particle and Nuclear Physics 63 (2009) 51–73



Neutrino distributions

$$n_{\nu_e}(\varepsilon_{\nu_e}) = \frac{96\varepsilon_{\nu_e}^2}{m_\mu^4} (m_\mu - 2\varepsilon_{\nu_e})$$

$$n_{\bar{\nu}_\mu}(\varepsilon_{\bar{\nu}_\mu}) = \frac{32\varepsilon_{\bar{\nu}_\mu}^2}{m_\mu^4} \left(\frac{3}{2}m_\mu - 2\varepsilon_{\bar{\nu}_\mu} \right)$$

the neutrino energy spectra from π^+ and μ^+ DAR. The vertical scale is arbitrary.

Reactor Neutrinos for low-energy interactions and neutrino oscillations

The LSND (Liquid Scintillator Reactor Neutrino Detector) was designed to search for the oscillation

$$\tilde{\nu}_{\mu} \rightarrow \tilde{\nu}_e$$

and obtained evidence for flavor change at the scale $\Delta m^2 \sim 0.2 \text{ eV}^2$

Although KARMEN did not confirm the LSND results, a combined analysis of the two experiments defines compatible regions:

W.C. Louis, Prog.Part. Nucl.Phys. 63 (09)51.

Both experiments used a **high-intensity 800 MeV proton beam** to produce a large number of π^+ pions which almost all decay (π^- and the produced μ^- are captured).

The BooNE neutrino detector (FNAL) is now also used for core collapse SN searches : [\[astro-ph/09103182v1\]](#)

Low-energy *Beta-beam neutrinos*

Recently, the use of *boosted β -decay radioactive nuclei* was proposed* to produce neutrino-beams of low energy (**beta-beam neutrinos**).

With this facility we obtain an intense, collimated and pure neutrino-beam **for searching Neutrino–Nucleus Interactions**.

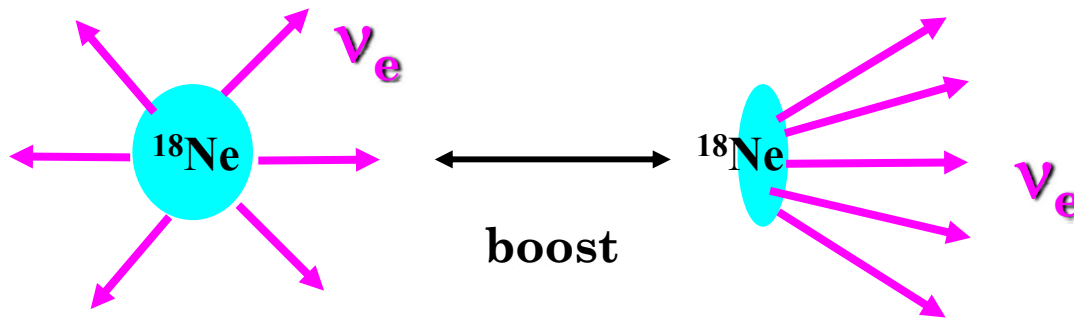
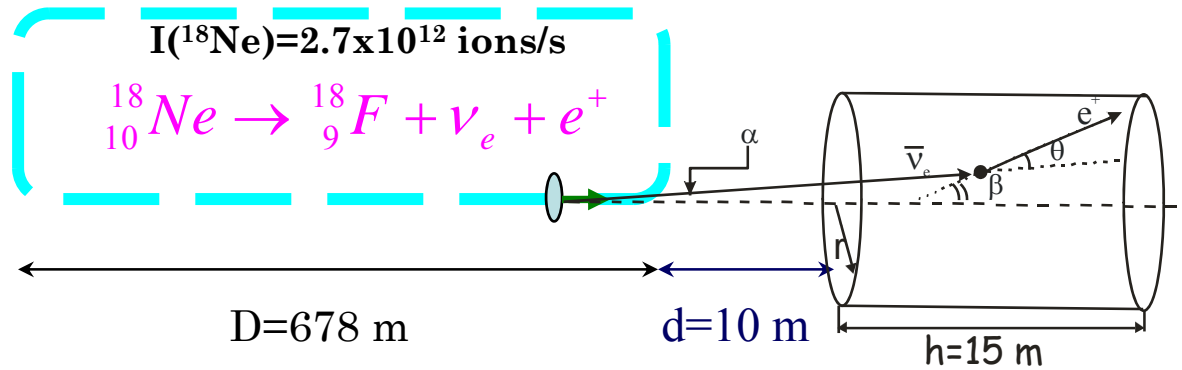
Such studies are important for various open issues in nuclear, particle physics, and astrophysics.

Low-energy Beta-beam neutrinos are useful for the interpretation of SN'v signals.

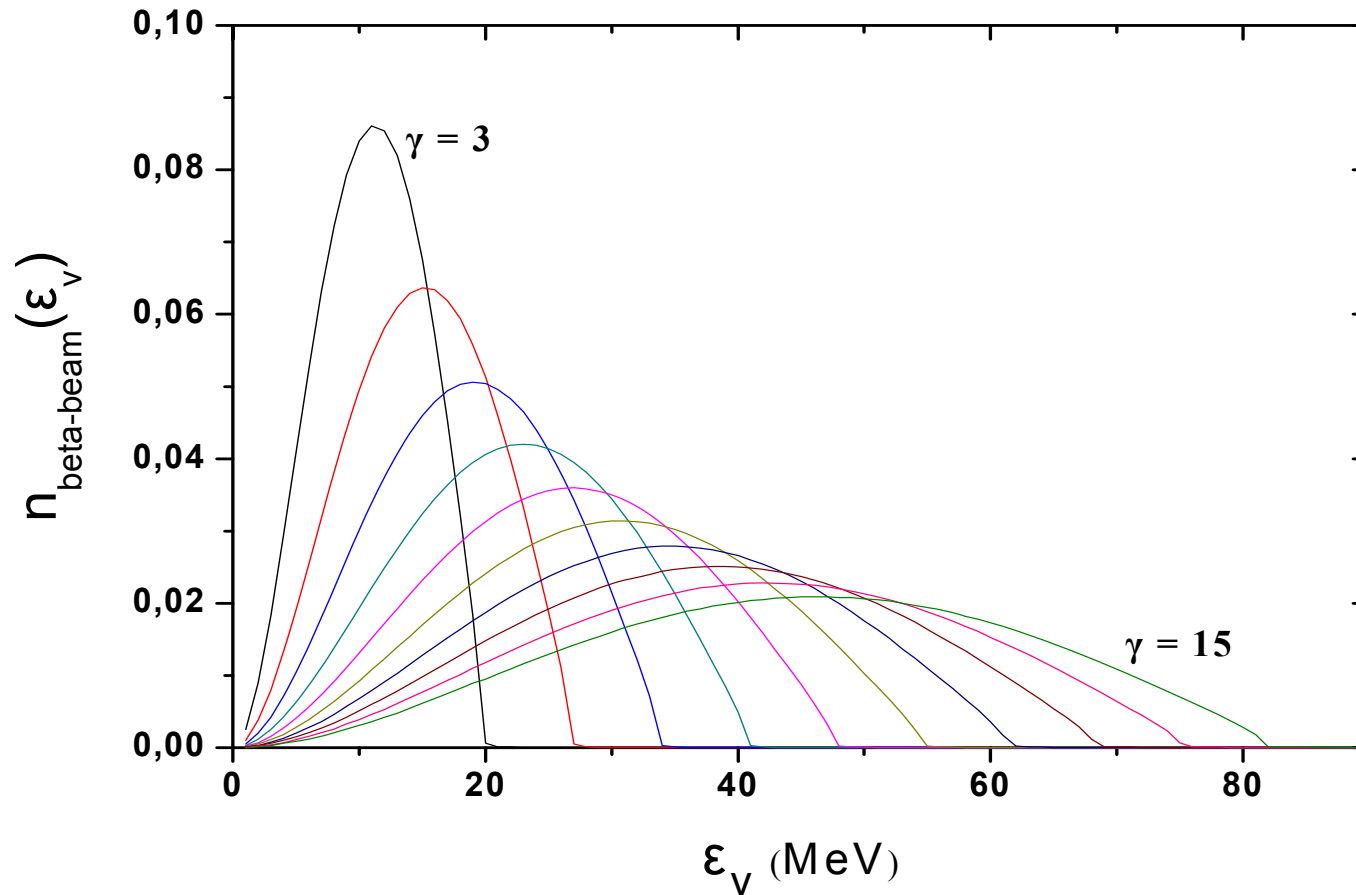
* C. Volpe, J. Phys. G, 30 (2004) L1–L6

Beta-beam neutrinos (radioactive sources)

Small Storage ring-detector



Energy-spectra of Boosted beta-beam neutrinos



For the interpretation of low-energy interactions (SN neutrinos, etc.) $\gamma < 15$
The distributions refer to ^{18}Ne but there are similar for ^6He too.

Solar Neutrinos

Neutrinos produced by the thermonuclear reactions in the Sun

Underground Neutrino detection Experiments.
(Homestake, GALLEX, SAGE, SK, SNO, etc.)

Measurements of solar neutrinos (KAMLAND, Borexino, etc.)
used to test the standard solar model (SSM).

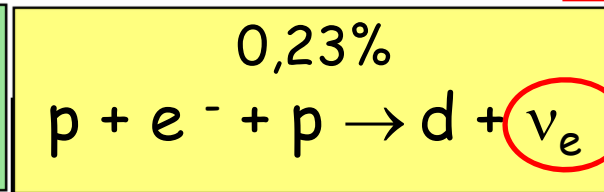
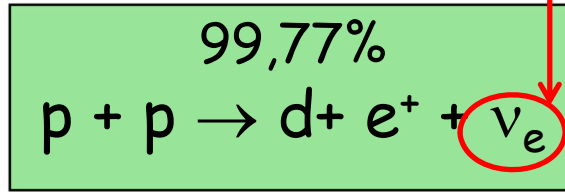
Recent probes, SNO+ experiment, aim to measure low-flux solar neutrinos
(*pep*, *CNO-cycle neutrinos*) to check the abundances of solar core and clarify if
the metallicity in the Sun is homogeneous

[K. Zuber, Inv. Talk, medex09].

The liquid scintillator detector of SNO+ experiment will be able to study low-energy solar neutrinos, geo-neutrinos, reactor neutrinos, supernova neutrino

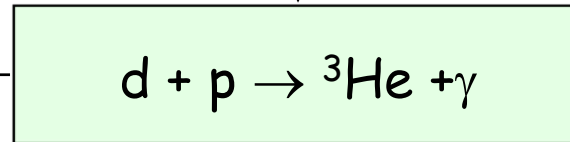
pp-chain reactions

$$Q_v \leq 0.420 \text{ MeV}$$



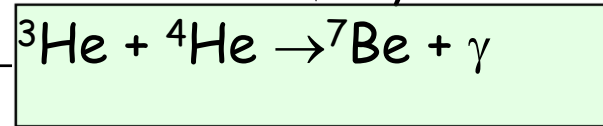
$$Q_v = 1.442 \text{ MeV}$$

84,7%

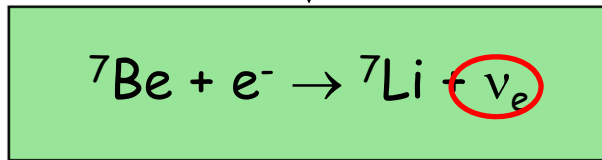


$\sim 2 \times 10^{-5} \%$

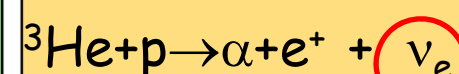
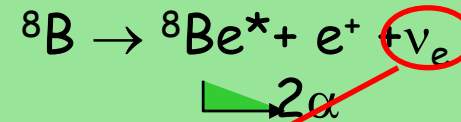
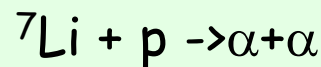
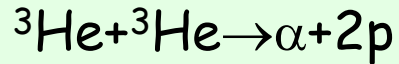
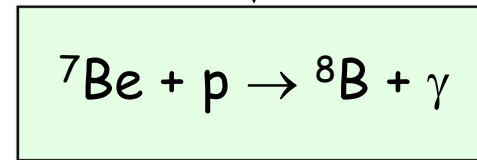
13,8%



13,78%



0,02%



pp I

pp II

$$Q_v < 15 \text{ MeV}$$

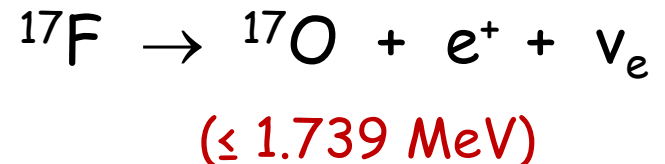
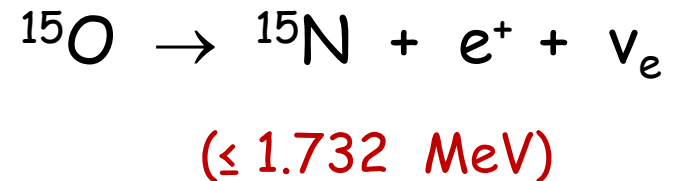
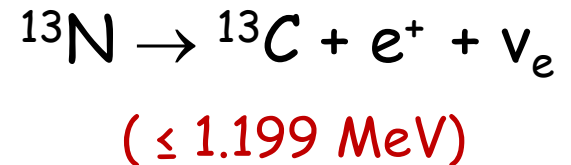
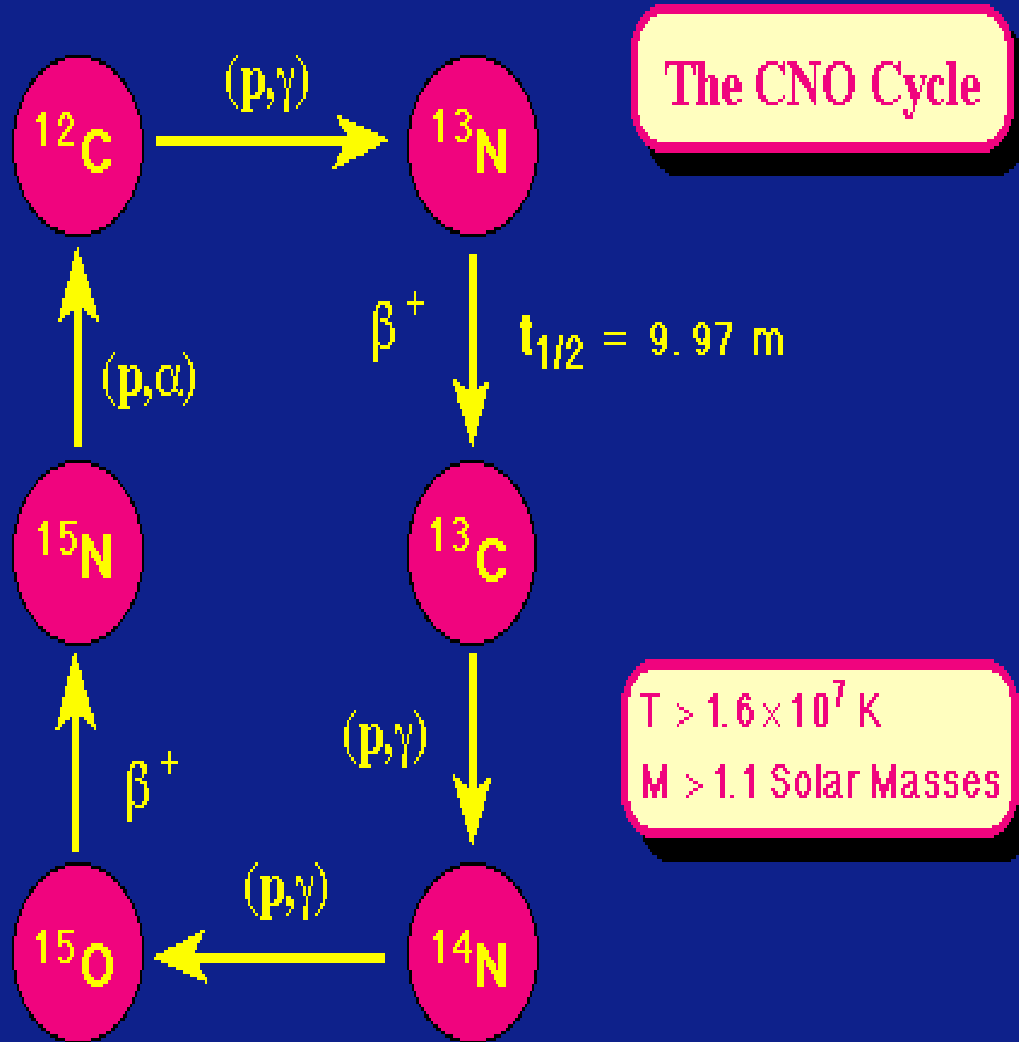
pp III

$$Q_v \leq 18.77 \text{ MeV}$$

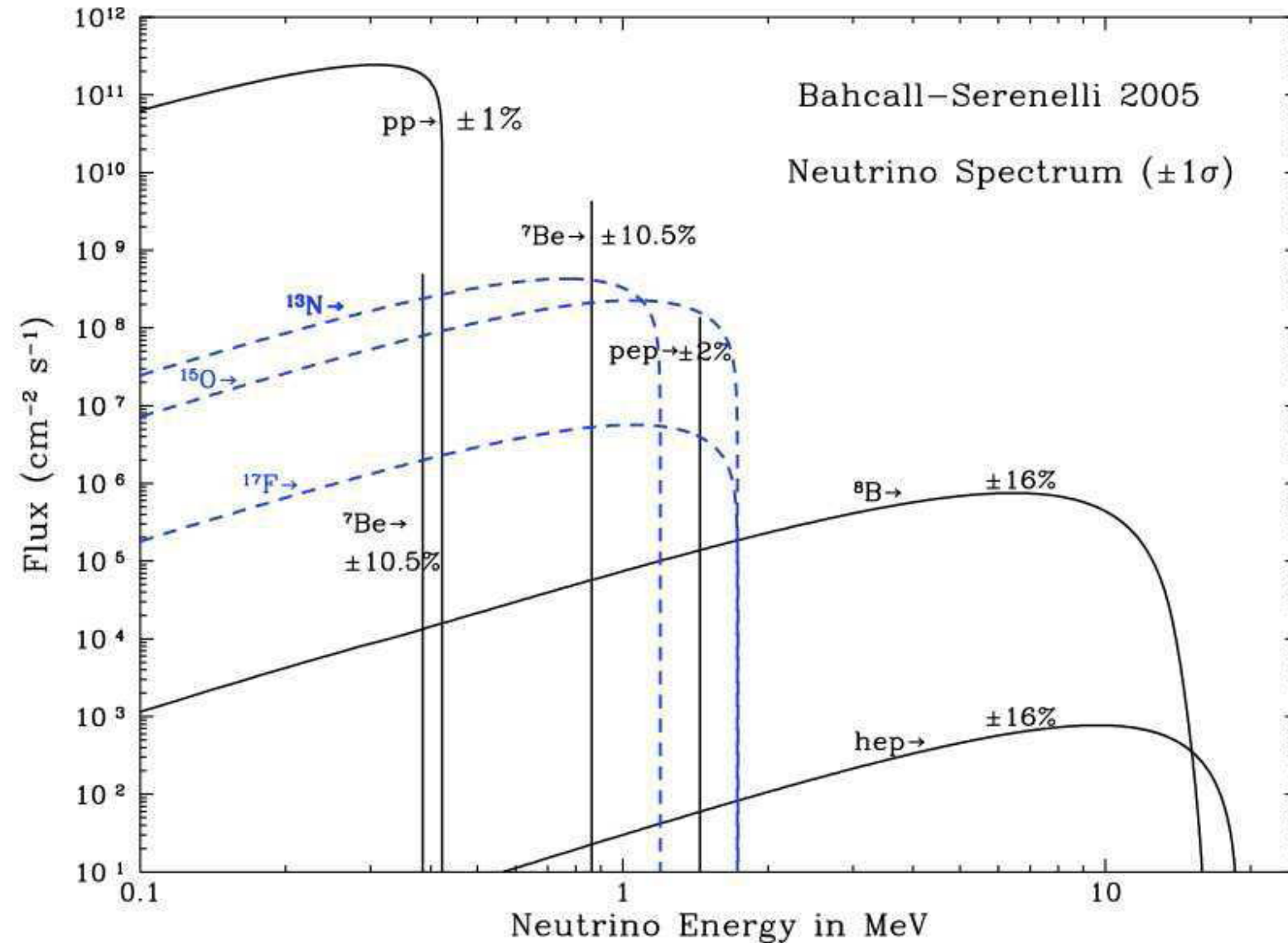
hep

The CNO cycle

2nd and 3rd generation stars like our Sun contain heavier elements (C,N,O,etc) (formed after the explosion, matter ejection of older stars)



Solar Neutrino fluxes (SSM)



Core-collapse Supernova Neutrinos

The energy-spectra of neutrinos emitted in a core collapse SN can be approximated as

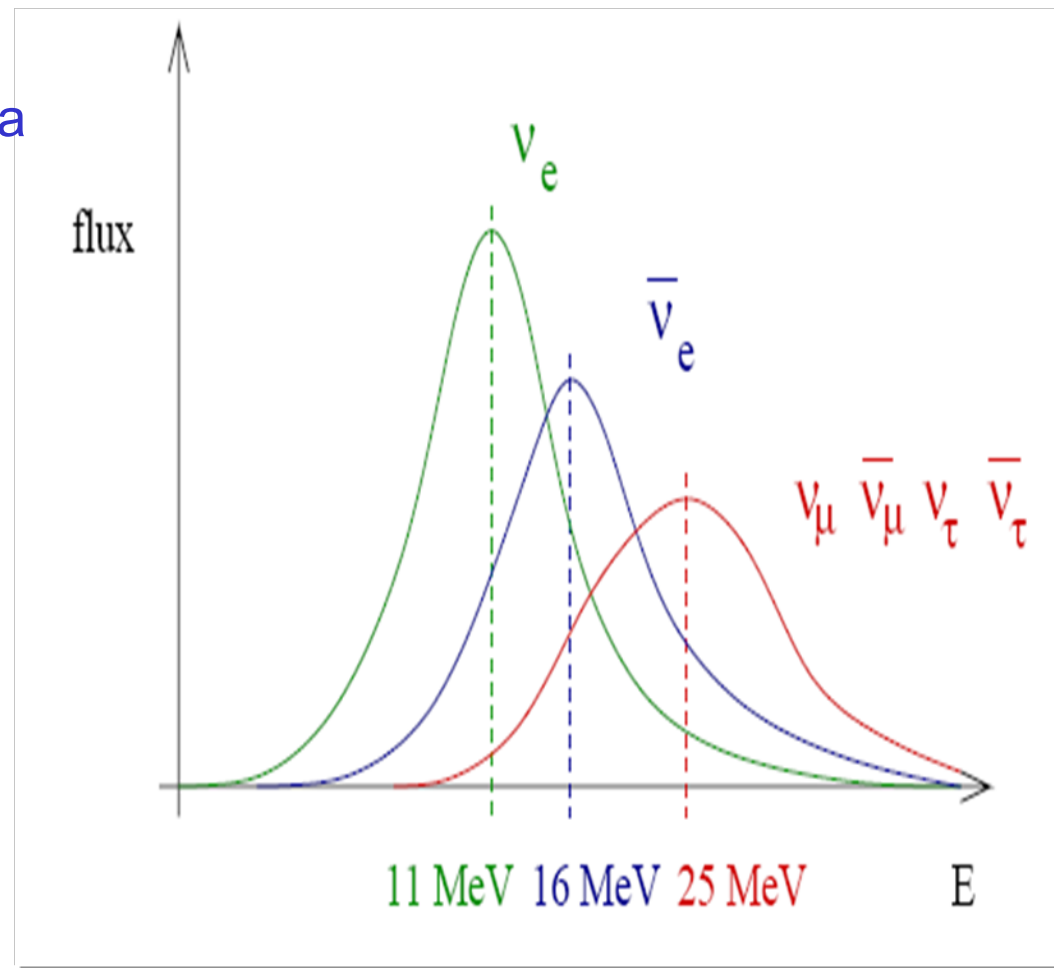
The Mean Energy of the SN spectra reflect the Temperature of the burning shell from where they escaped

$$\langle E_{\nu_e} \rangle \approx 11 \text{ MeV}$$

$$\langle E_{\bar{\nu}_e} \rangle \approx 16 \text{ MeV}$$

$$\langle E_{\nu_{\mu\tau}} \rangle \approx 25 \text{ MeV}$$

SN Energy spectra



Parameterization of SN- ν energy-spectra

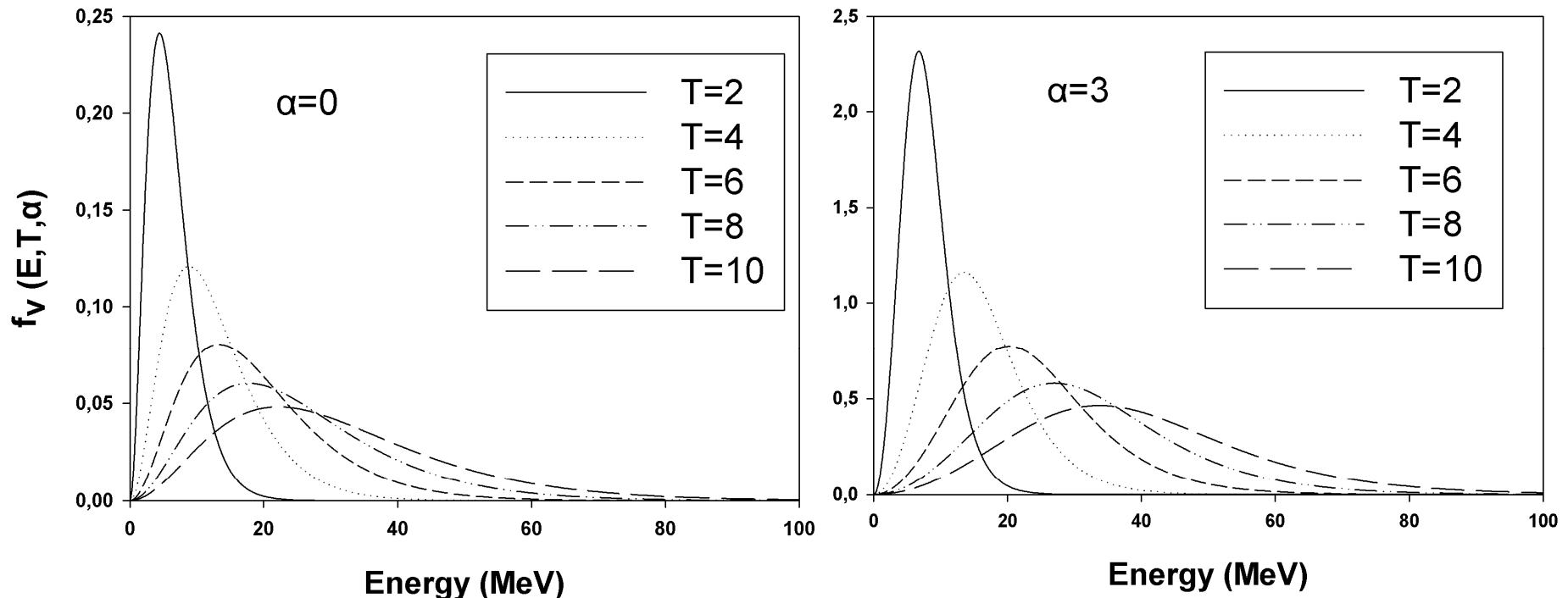
In modeling core-collapse SN and studies of nuclear response to SN- ν , various parametrizations are used :

$$f(E_\nu) = \frac{1}{F_2(\alpha)T^3} \frac{E_\nu^2}{\exp[(E_\nu/T) - \alpha] + 1}$$

Fermi-Dirac
distribution

T = Neutrino Temperature α = Chemical Potential

Supernova neutrino spectra



Power-Law distribution for SN- ν

The SN- ν energy-spectra can be accurately parameterized with a power-law distribution of the form

$$PL(\langle \varepsilon_\nu \rangle, \varepsilon_\nu, \alpha) = \left(\frac{\varepsilon_\nu}{\langle \varepsilon_\nu \rangle} \right)^\alpha e^{-(a+1) \frac{\varepsilon_\nu}{\langle \varepsilon_\nu \rangle}}$$

$\langle \varepsilon_\nu \rangle$ = the average neutrino energy

α = spectral pinching parameter

These parameters allow to adjust the width w

$$w = \sqrt{\langle \varepsilon_\nu^2 \rangle - \langle \varepsilon_\nu \rangle^2}$$

Beta-beam neutrino Energy-spectra for SN- ν

The low-energy beta-beam neutrinos can be exploited to interpret a SN- ν signal

Linear combinations of normalized beta-beam spectra are used

$$n_{N_\gamma}(\varepsilon_\nu) = \sum_{i=1}^N \alpha_i n_{\gamma_i}(\varepsilon_\nu)$$

The fitting of coefficients α_i and boost factors γ_i to a SN-neutrino spectrum is achieved by minimizing the integral

$$\int_{\varepsilon_\nu} d\varepsilon_\nu \left| n_{N_\gamma}(\varepsilon_\nu) - n_{SN}(\varepsilon_\nu) \right|$$

Neutrino-nucleus interactions

During the last decades, neutrino-nucleus interactions, β -decay modes, nuclear μ -capture, etc., have deepen our knowledge on :

- 1) the fundamental electro-weak interactions
- 2) our understanding on nuclear structure

Such precious information inspired significant probes within and beyond the SM and provided valuable interpretations to experiments looking for:

- 1) neutrino detection probes
- 2) Neutrino oscillation probes
- 3) neutrino masses

Neutrino–nucleus reactions play important role in astrophysical processes (SN nucleo-synthesis, etc.).

Uncertainties on astrophysical interactions of neutrinos raised numerous questions, still unanswered.

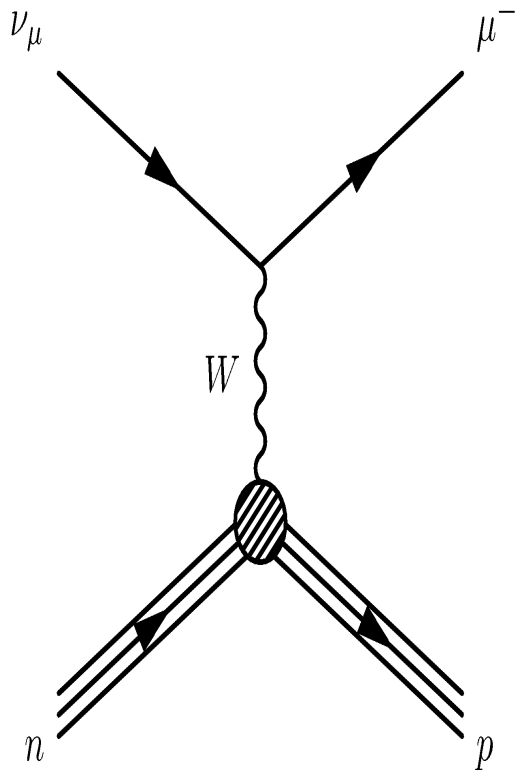
Our knowledge is limited by the understanding of neutrino-nucleus cross sections based on reliable interaction Hamiltonians and w-f.

Theory of Neutrino-nucleus interactions

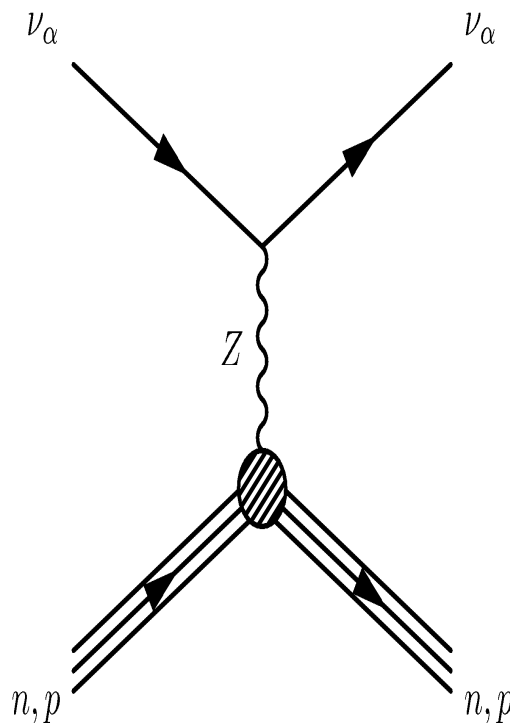
NC: Neutral current interaction (mediated by the Z-boson)

CC: Charged-currents interaction (mediated by W^+ or W^- boson)

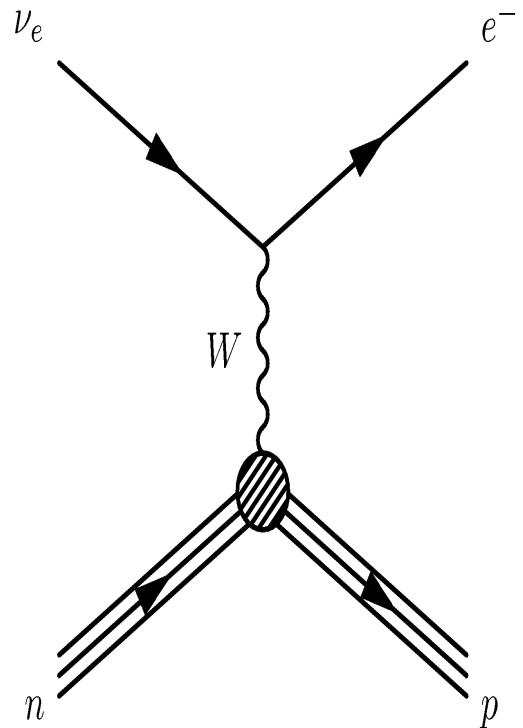
ν_μ CC Event



NC Event



ν_e CC Event



Neutrino-Nucleus reactions

There are **four types** of ν -nucleus reactions for ν -detection

- Charged-current reactions ($l = e, \mu, \tau$)

$$\nu_l + {}_Z A_N \longrightarrow {}_{Z+1} A_{N-1}^* + l^-$$

$$\bar{\nu}_l + {}_Z A_N \longrightarrow {}_{Z-1} A_{N+1}^* + l^+ .$$

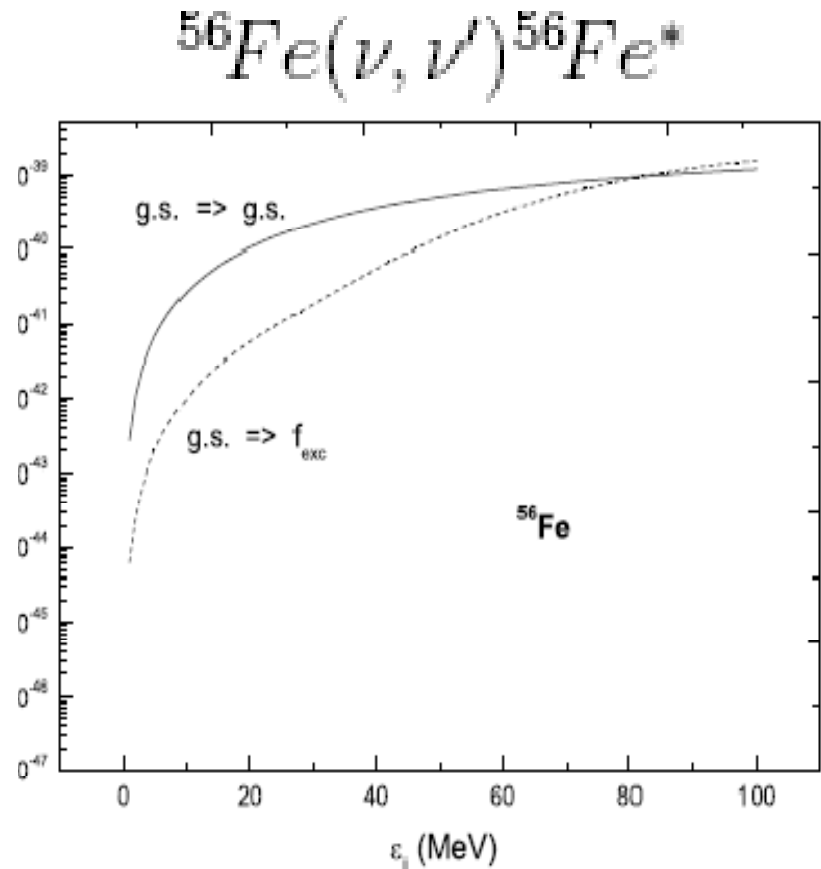
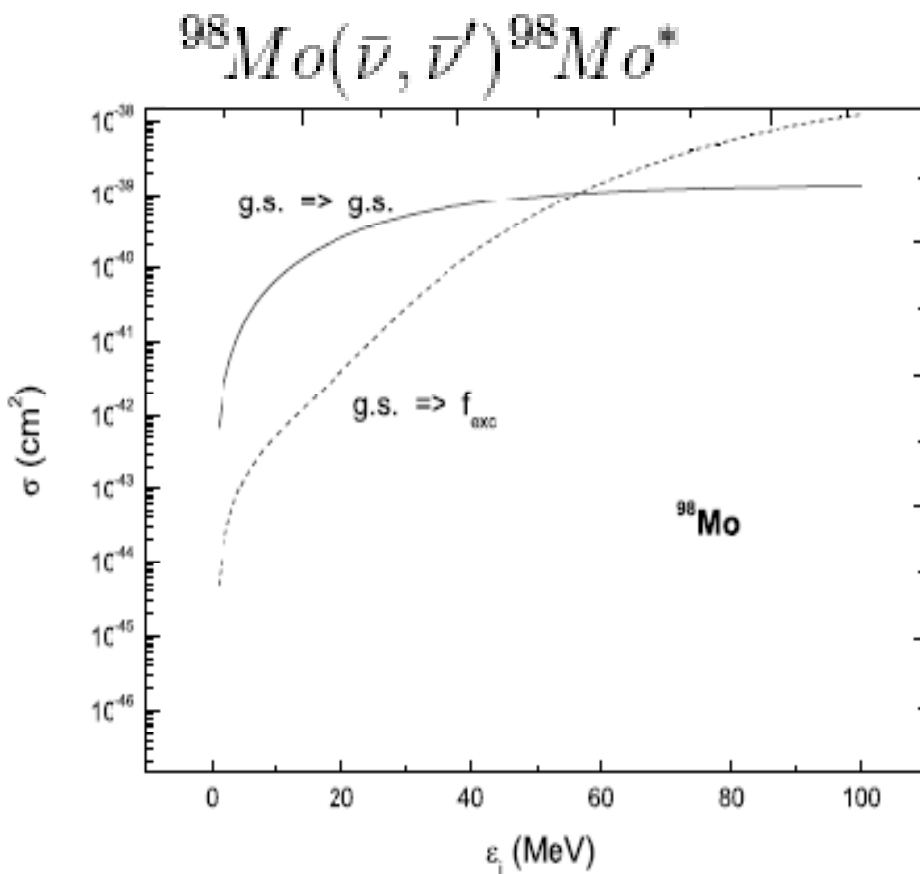
- Neutral-current reactions (Coherent Channel Possible)

$$\nu + {}_Z A_N \longrightarrow {}_Z A_N^* + \nu'$$

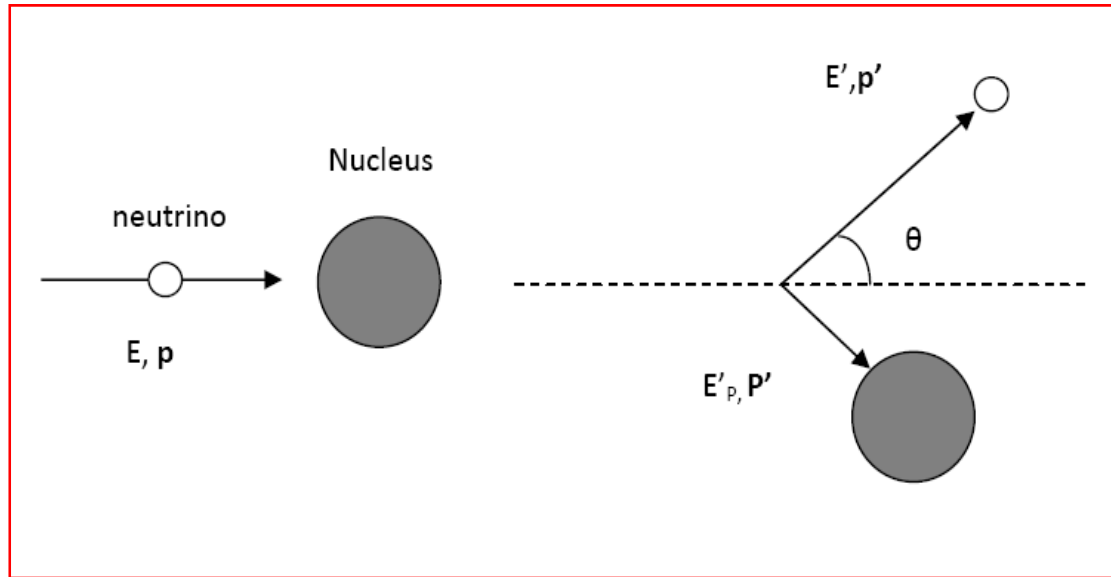
$$\bar{\nu} + {}_Z A_N \longrightarrow {}_Z A_N^* + \bar{\nu}'$$

Coherent ν -Nucleus processes

For low-energy neutrinos (solar neutrinos), the coherent channel dominates the total cross section ([tsk, et al., NPA, submitted](#)).



Coherent ν -Nucleus scattering (Neutral-Currents)



- The **Coherent** differential cross section reads

$$\frac{d\sigma}{d(\cos\Phi)} = G^2 \frac{\sin^2\theta_w}{2\pi} A^2 E_\nu^2 (1 + \cos\Phi)$$

Φ = scattering angle, θ_w = Weinberg angle, E_ν = incoming-neutrino energy

- Nucleons contribute coherently (coherent process dominates)

Nuclear recoil

In experimental studies of Coherent neutrino-nucleus scattering, the only observable is the average recoil energy

$$\langle E_N \rangle = \frac{2}{3A} \left(\frac{E_\nu}{1 \text{ MeV}} \right)^2 \text{ keV}$$

- (i) Proportional to E_ν^2 , inversely proportional to A (mass number)
- (ii) Including other experimental criteria, ^{28}Si , ^{32}S best choices

The calculations start from the quantity $d\sigma/dT_A$

$$\left(\frac{d\sigma}{dT_A} \right)_{\text{weak}} = \frac{G_F^2 A m_N}{2\pi} (N^2/4) F_{\text{coh}}(T_A, E_\nu)$$

$$F_{\text{coh}}(T_A, E_\nu) = F^2(q^2) \left(1 + \left(1 - \frac{T_A}{E_\nu} \right)^2 \right) - \frac{A m_N T_A}{E_\nu^2}$$

$$q^2 = T_A^2 + 2A m_N T_A$$

Coherent neutrino-nucleus scattering

with

$$F(Q^2) = \frac{1}{Q_w} [N F_n(Q^2) - Z(1 - 4 \sin^2 \theta_w) F_p(Q^2)]$$

The quantity Q_w is

$$Q_w = N - Z(1 - 4 \sin^2 \theta_w)$$

Finally one needs the neutron density of the target

$$F(Q^2) = \frac{1}{Q_w} \int [\rho_n(r) - (1 - 4 \sin^2 \theta_w) \rho_p(r)] \frac{\sin(Qr)}{Qr} r^2 dr$$

Coherent neutrino-Nucleus Scattering probes

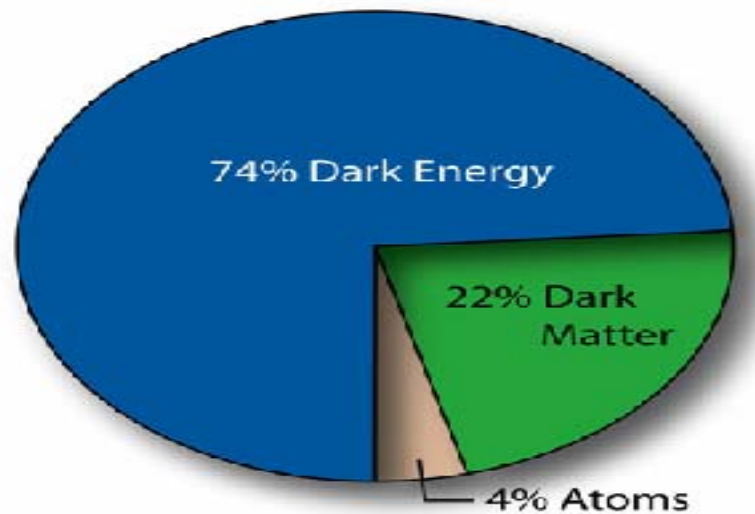
Absorber	E_{Thresh} (eV)	Rate ($K g^{-1} d^{-1}$)
<i>S1</i>	100	40
	300	22
	500	14
<i>Al₂O₃</i>	100	36
	300	22
	500	14
<i>Ge</i>	100	85
	300	24
	500	8
<i>Pb</i>	50	201
	100	83
	300	4
	500	< 1



Such studies could be done in conjunction with direct detection of Cold Dark Matter (cryogenic detectors).

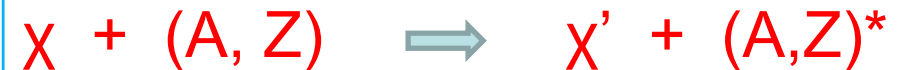
Counting rates for Coherent Neutrino-Nucleus Scattering.
The Reactor-Neutrino flux is $10^{13} \text{ cm}^{-2} \text{ sec}^{-1}$

Direct CDM detection: CDM-particle Scattering off nuclei



The Content of the universe:
Dark Energy $\approx 73\%$, Atoms $\approx 4\%$
Cold Dark Matter $\approx 23\%$

LSP-nucleus scattering



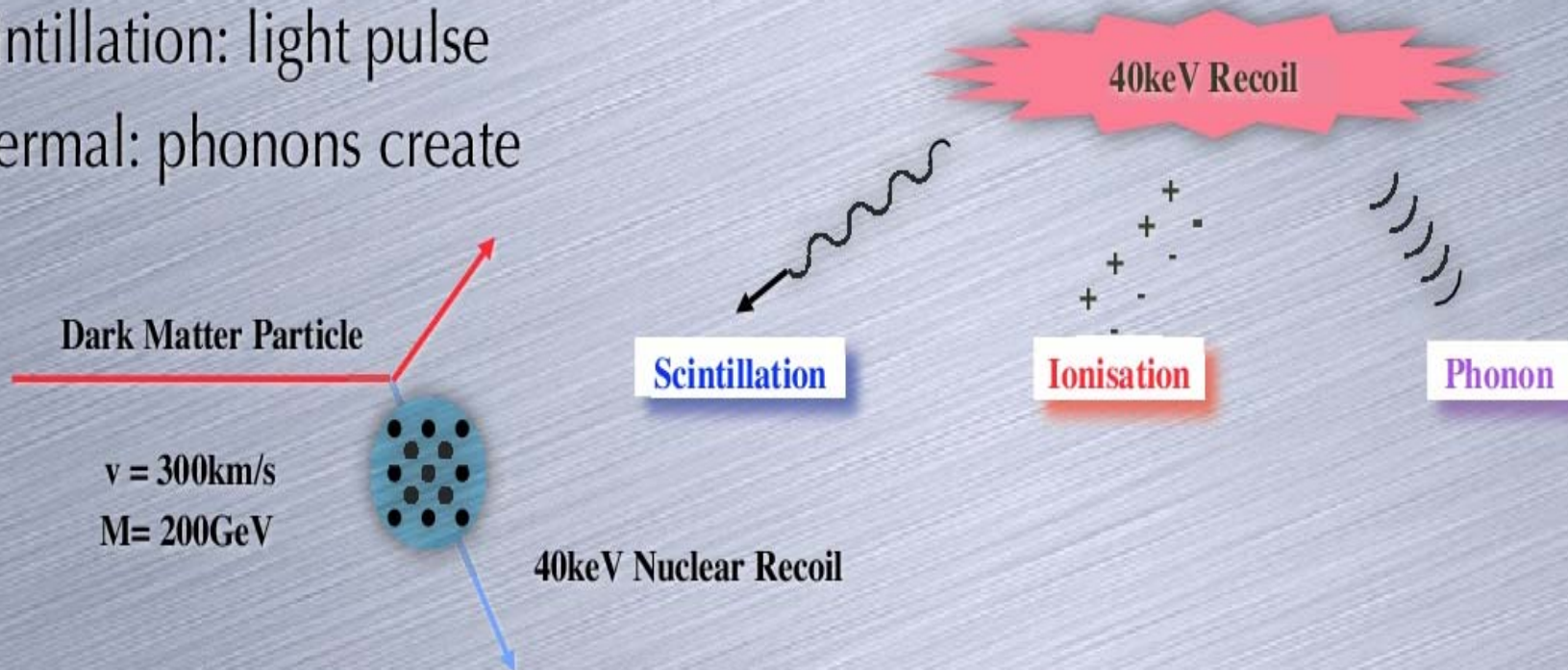
A) Coherent process is possible : Vector & Axial-Vector Currents

B) Dominance of Axial-Vector contributions

(Odd-A nuclear targets : ^{73}Ge , ^{127}I , ^{115}In , $^{129,131}\text{Xe}$)

Signals of CNN scattering and direct detection of WIMPs

- Ionisation: free electric charge release
- Scintillation: light pulse
- Thermal: phonons create



The energy transfer is very low (30-40 keV)

Inelastic ν -Nucleus scattering (Cross section)

In Walecka-Donnelly method [PRC 6 (1972)719, NPA 201(1973)81]

$$\frac{d^2\sigma_{i\rightarrow f}}{d\Omega d\omega} = \frac{G^2}{\pi} \frac{\varepsilon_f^2}{(2J_i + 1)} \left(\sum_{J=0}^{\infty} \sigma_{CL}^J + \sum_{J=1}^{\infty} \sigma_T^J \right)$$

where

$$\omega = \varepsilon_i - \varepsilon_f \quad q = |\mathbf{q}| = [\omega^2 + 2\varepsilon_i\varepsilon_f (1 - \cos \Phi)]^{\frac{1}{2}}$$

1st sum = Coulomb-Longitudinal contr.

2nd sum = Transverse contr.

One of the goals is: To investigate the Response as low-energy neutrino detectors of :

- i) Te, Cd, Zn –isotopes (COBRA $\beta\beta$ -decay experiment)
- ii) Mo-isotopes (MOON $\beta\beta$ -decay experiment)

- i) K.Zuber, Phys. Lett. B 571(03)148, talk at MEDEX-09, June 15-19, 2009, Prague.
- ii) H. Ejiri, Phys. Rep. 338 (2000)265; H. Ejiri et al., Phys.Lett. B 530 (02)265;

Compact expressions for the 7-basic reduced ME

For H.O. basis, all basic reduced ME take analytic compact forms

$$\langle j_1 || T^J || j_2 \rangle = e^{-y} y^{\beta/2} \Pi(y) = e^{-y} y^{\beta/2} \sum_{\mu=0}^{n_{max}} \mathcal{P}_{\mu}^J y^{\mu}.$$

The Polynomials of even terms in q have constant coefficients as

$$\Pi(y) = \sum_{\mu=0}^{n_{max}} \mathcal{P}_{\mu}^J y^{\mu}, \quad y = \frac{q^2 b^2}{4}$$

$$n_{max} = (N_1 + N_2 - \beta)/2.$$

V.Chasioti, TSK, Czech. J. Phys. 52 (2002)467; Nucl. Phys. A, 829 (2009) 234

Advantages of the above formalism (FORTRAN Code) :

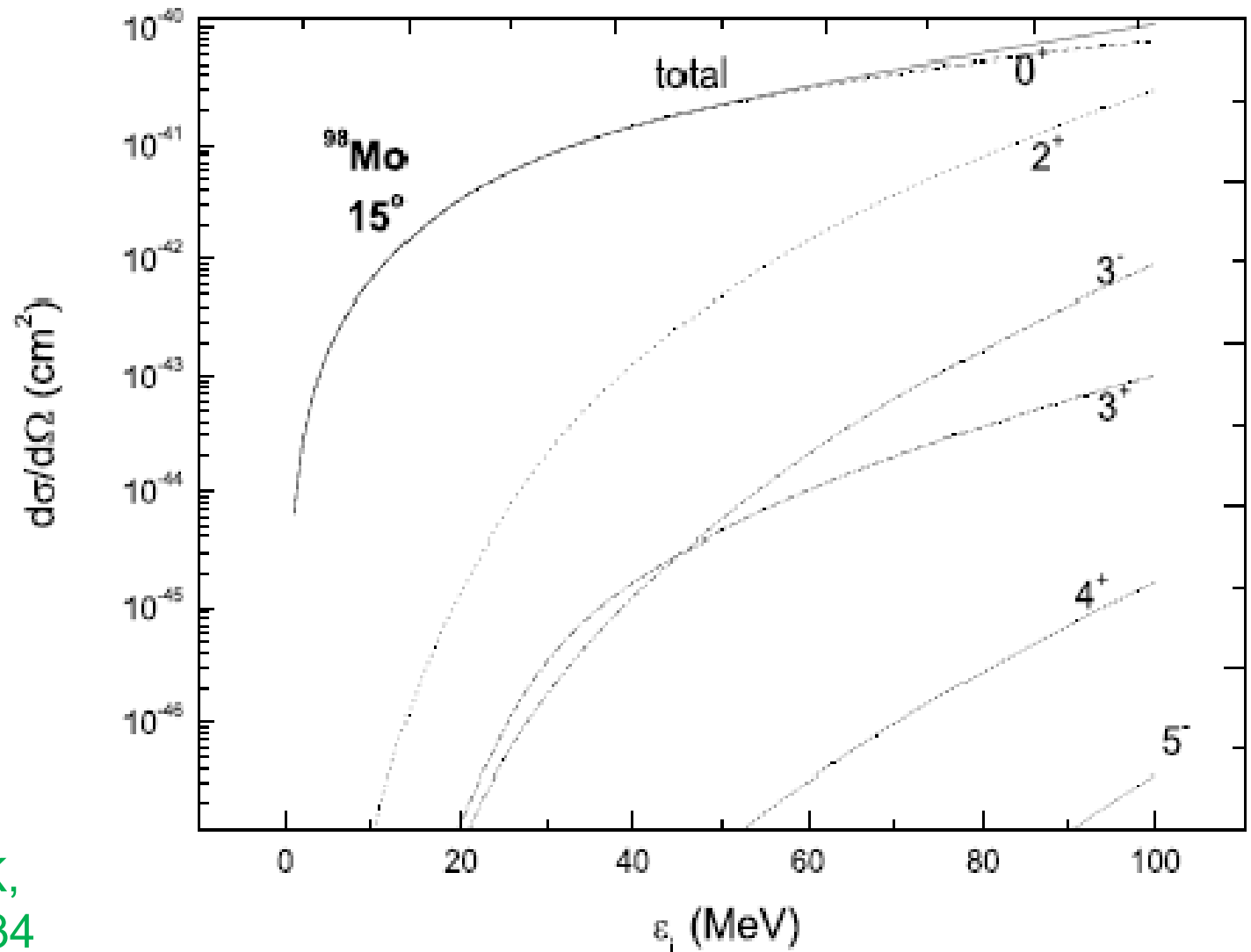
- (i) The coefficients \mathcal{P}^J are calculated once (reduction of computer time)
- (ii) They can be used for phenomenological description of ME
- (iii) They are useful for other bases sets (expansion in HO wave-functions)

Neutrino-nucleus cross sections

- A). On the basis of the original calculations we could study:
- i) State-by-state calculations of $d\sigma/d\Omega$
 - ii) The contribution of the individual multipolarities
 - iii) Dominance of Axial-Vector contributions in σ
- B). Nuclear detector response to ν -sources:
- i) solar neutrino spectra
 - ii) SN neutrino spectra
 - iii) reactor neutrino spectra

State-by-state calculations of $d\sigma/d\Omega$

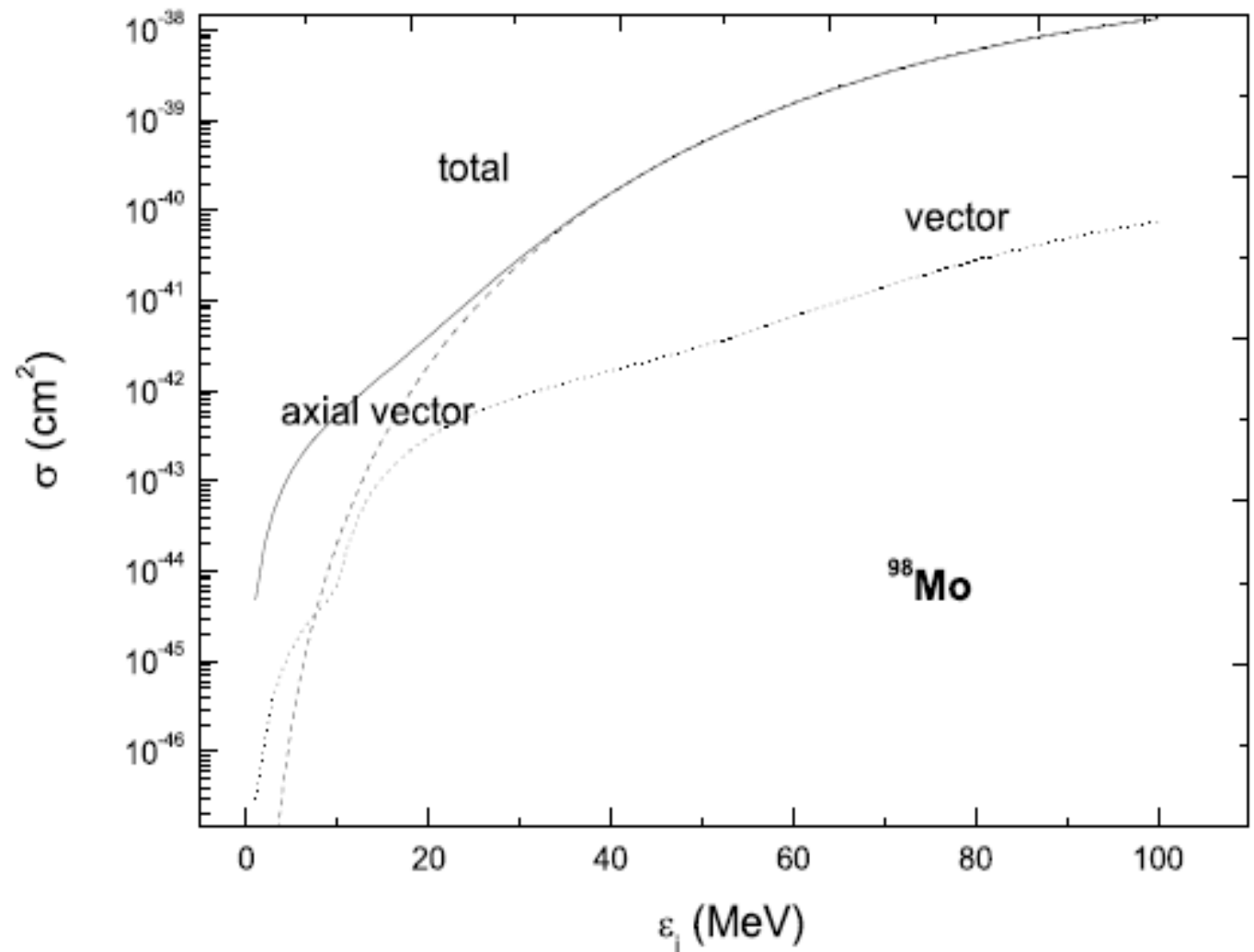
$${}^{98}\text{Mo}(\bar{\nu}, \bar{\nu}') {}^{98}\text{Mo}^*$$



Dominance of Axial-Vector contributions in σ

$${}^{98}\text{Mo}(\bar{\nu}, \bar{\nu}') {}^{98}\text{Mo}^*$$

${}^{98}\text{Mo}$



Summary and Outlook

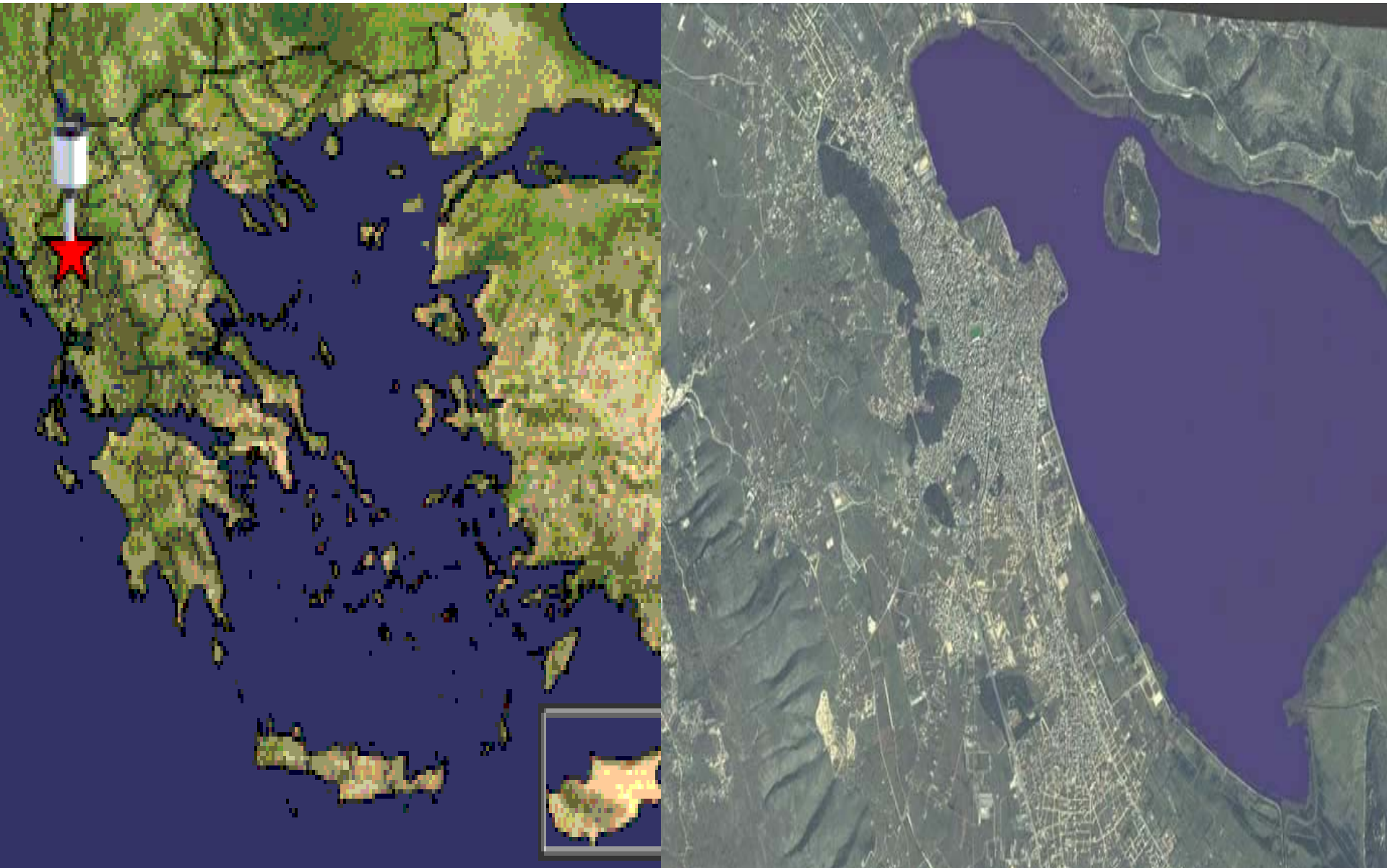
(i) Low-energy neutrinos searches

- i. Solar, SN neutrinos
- ii. Reactor and Beta-beam neutrino

Require: *a). Reliable neutrino-nucleus reaction cross sections and*
b). The Response Nuclear ν -detectors to various ν -spectra

(ii) Recent and future nuclear-recoil experiments, with very low-threshold energy, could measure Coherent Neutrino-Nucleus Scattering in conjunction with experiments of Direct detection of CDM particles and 0ν -double-beta-decay.

Ioannina, the city of Silver and Gold



Thank you for your attention